

Breaking internal waves and ocean diapycnal diffusivity: A case study based upon dynamically downscaling of high-resolution global ocean model simulations into a regional domain

*Kayhan Momeni*¹

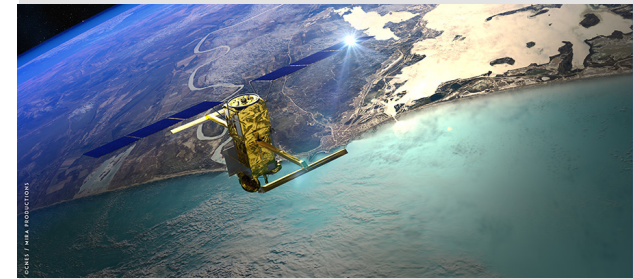
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1. Numerical Setup
2. Theoretical Analysis
3. Shear-induced mixing

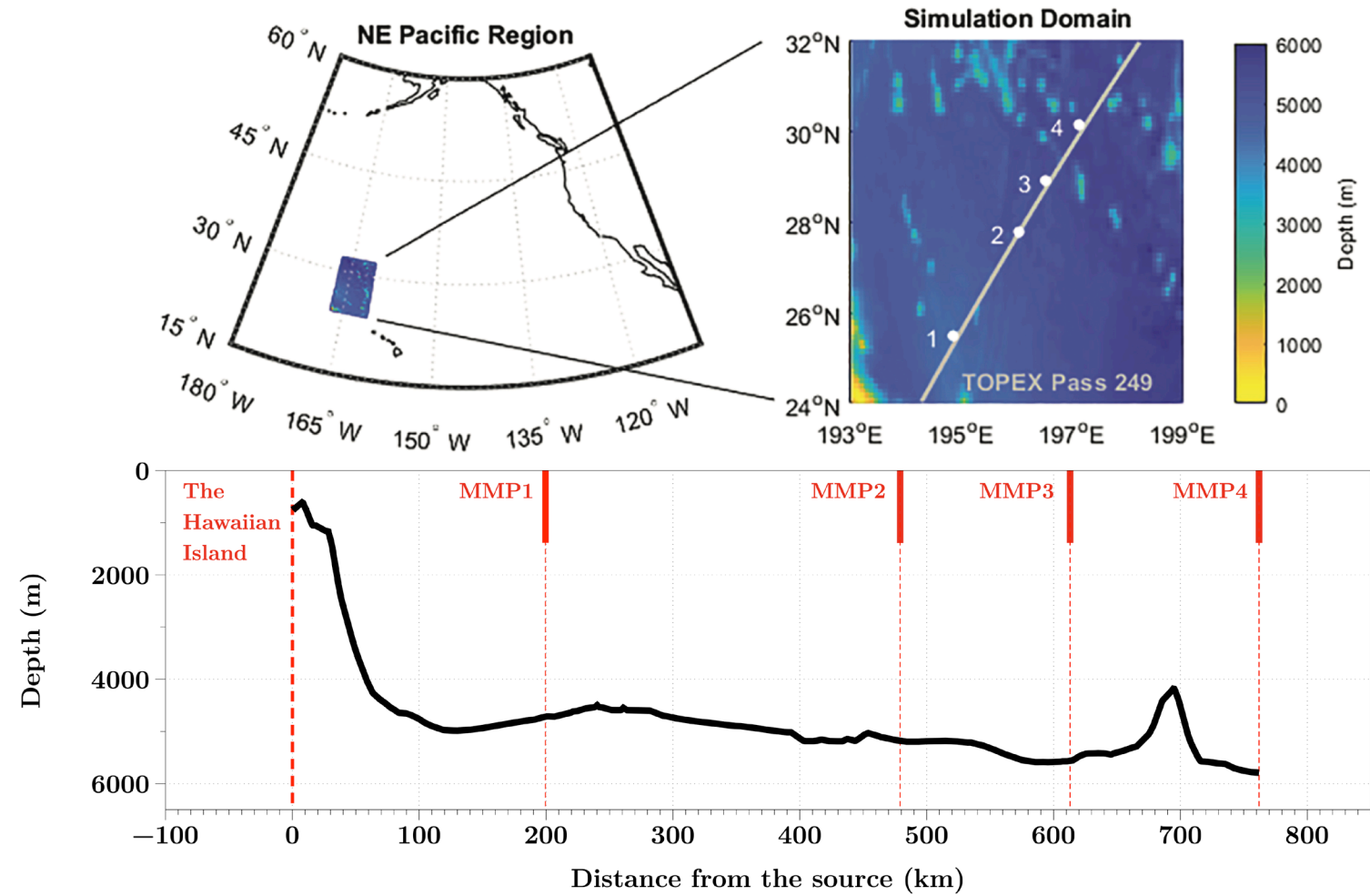


Image source (top): Nelson AD, Arbic BK, Menemenlis D, Peltier WR, Alford MH, Grisouard N, Klymak JM. Improved internal wave spectral continuum in a regional ocean model. *Journal of Geophysical Research: Oceans*. 2020 May;125(5):e2019JC015974.



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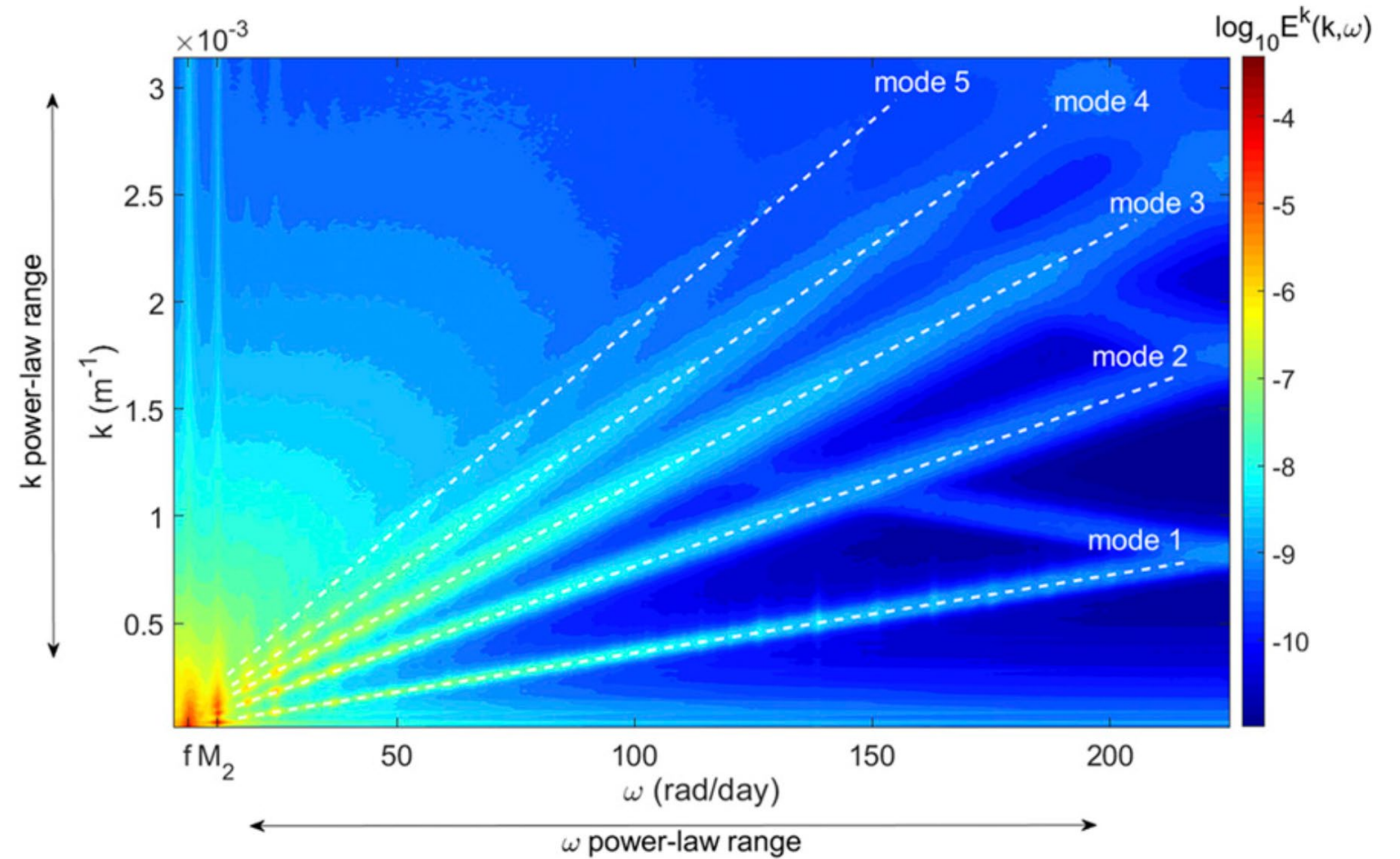


Image source: Pan Y, Arbic BK, Nelson AD, Menemenlis D, Peltier WR, Xu W, Li Y. Numerical investigation of mechanisms underlying oceanic internal gravity wave power-law spectra. *Journal of Physical Oceanography*. 2020 Sep 1;50(9):2713-33.



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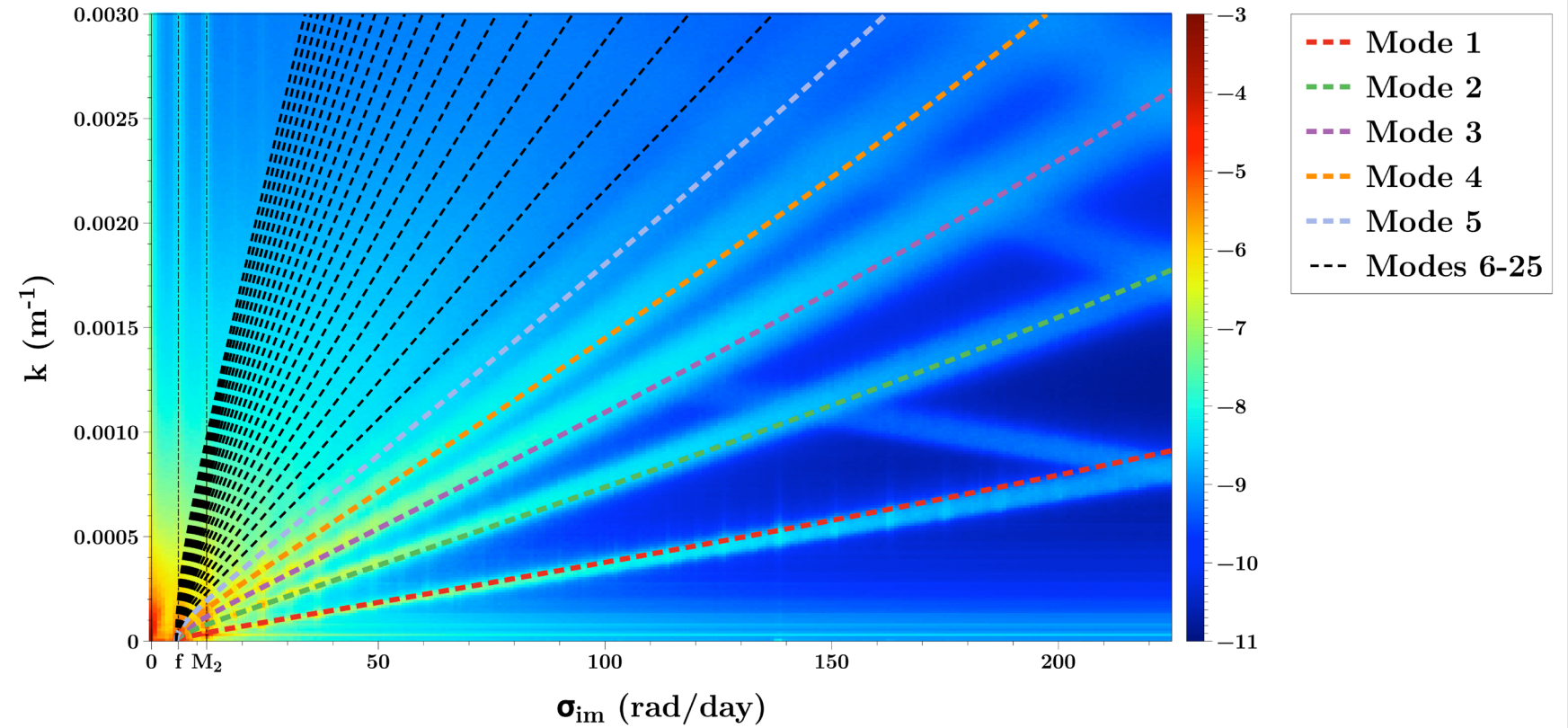
$$\left\{ \begin{aligned}
 U_x + V_y + W_z &= 0 \\
 U_t + UU_x + VU_y + WU_z &= -\frac{P_x^{eff}}{\rho_0} + \nu^{eff} \nabla^2 U + \frac{\partial \nu^{eff}}{\partial z} (U_z + W_x) + fV \\
 V_t + UV_x + VV_y + WV_z &= -\frac{P_y^{eff}}{\rho_0} + \nu^{eff} \nabla^2 V + \frac{\partial \nu^{eff}}{\partial z} (V_z + W_y) - fU \\
 W_t + UW_x + VW_y + WW_z &= -\frac{P_z^{eff}}{\rho_0} + \nu^{eff} \nabla^2 W + 2\frac{\partial \nu^{eff}}{\partial z} W_z - g\frac{\rho}{\rho_0} \\
 \rho_t + U\rho_x + V\rho_y + W\rho_z &= \kappa^{eff} \nabla^2 \rho + \frac{\partial \kappa^{eff}}{\partial z} \rho_z
 \end{aligned} \right.$$

$$\left\{ \begin{aligned}
 \{U', V', P'\}(x, y, z, t) &= \left[\sum_{j=1}^J \{U_j, V_j, P_j\} \cos\left(\frac{j\pi z}{L}\right) \right] e^{(\sigma t - i(kx + ly))}, \\
 \{W', \rho'\}(x, y, z, t) &= \left[\sum_{j=1}^J \{W_j, R_j\} \sin\left(\frac{j\pi z}{L}\right) \right] e^{(\sigma t - i(kx + ly))}.
 \end{aligned} \right.$$

$$\left\{ \begin{aligned}
 \sigma \Omega_j &= -\nu_{mj}^{cc} k_m^2 \Omega_m - \nu_{zmj}^{cs} \left(\frac{m\pi}{L}\right) \Omega_m - f \left(\frac{j\pi}{L}\right) W_j \\
 \sigma k_j^2 W_j &= \nu_{zmj}^{sc} k_m^2 \left(\frac{m\pi}{L}\right) W_m - \nu_{mj}^{ss} k_m^2 \left(\frac{m^2 \pi^2}{L^2}\right) W_m + \nu_{zzmj}^{ss} \left(\frac{m^2 \pi^2}{L^2}\right) W_m \\
 &\quad + \nu_{zmj}^{sc} \left(\frac{m^3 \pi^3}{L^3}\right) W_m - \nu_{zmj}^{sc} (k^2 + l^2) \left(\frac{m\pi}{L}\right) W_m - \nu_{zzmj}^{ss} (k^2 + l^2) W_m \\
 &\quad - \nu_{mj}^{ss} k_m^2 (k^2 + l^2) W_m + 2\nu_{zmj}^{sc} (k^2 + l^2) \left(\frac{m\pi}{L}\right) W_m - \frac{g}{\rho_0} (k^2 + l^2) R_j \\
 &\quad + f \left(\frac{j\pi}{L}\right) \\
 \sigma R_j &= -\rho_{zmj}^{ss} W_m - k_m^2 \kappa_{mj}^{ss} R_m + \kappa_{zmj}^{sc} \left(\frac{m\pi}{L}\right) R_m
 \end{aligned} \right.$$

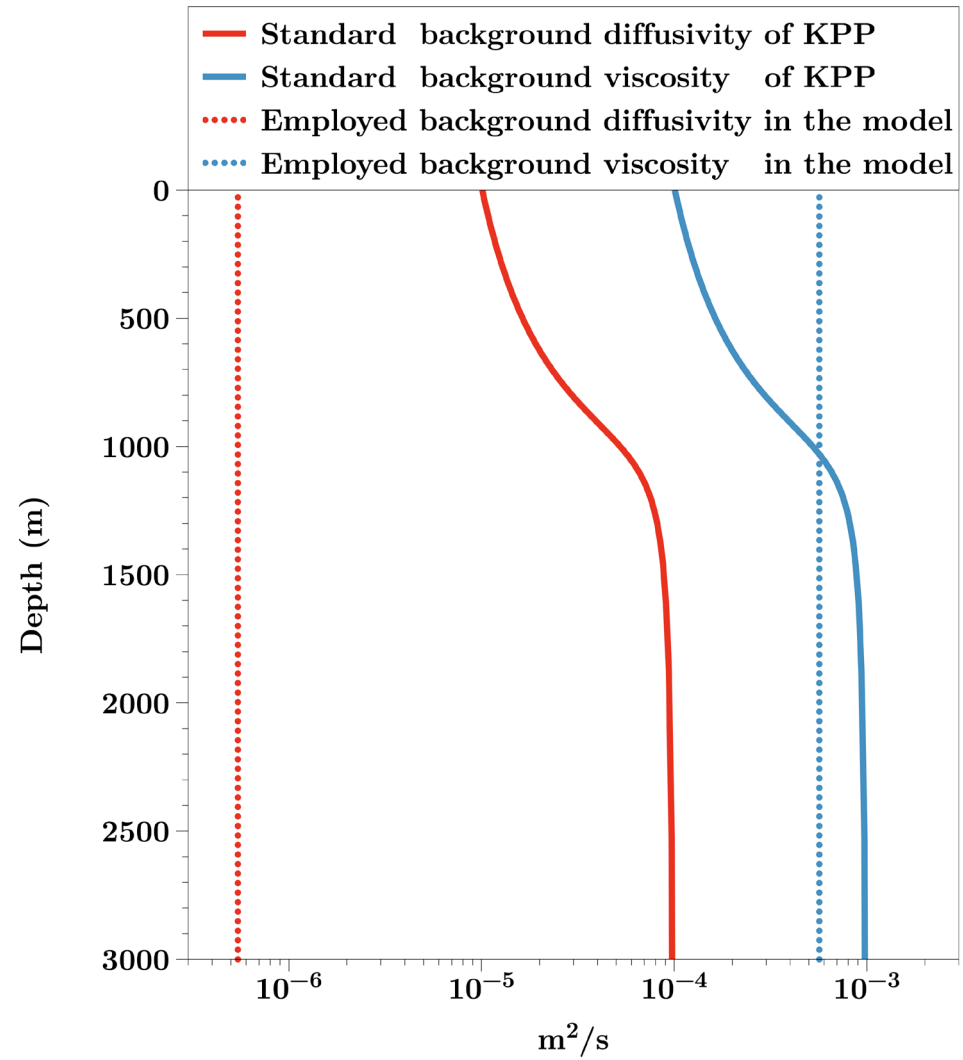


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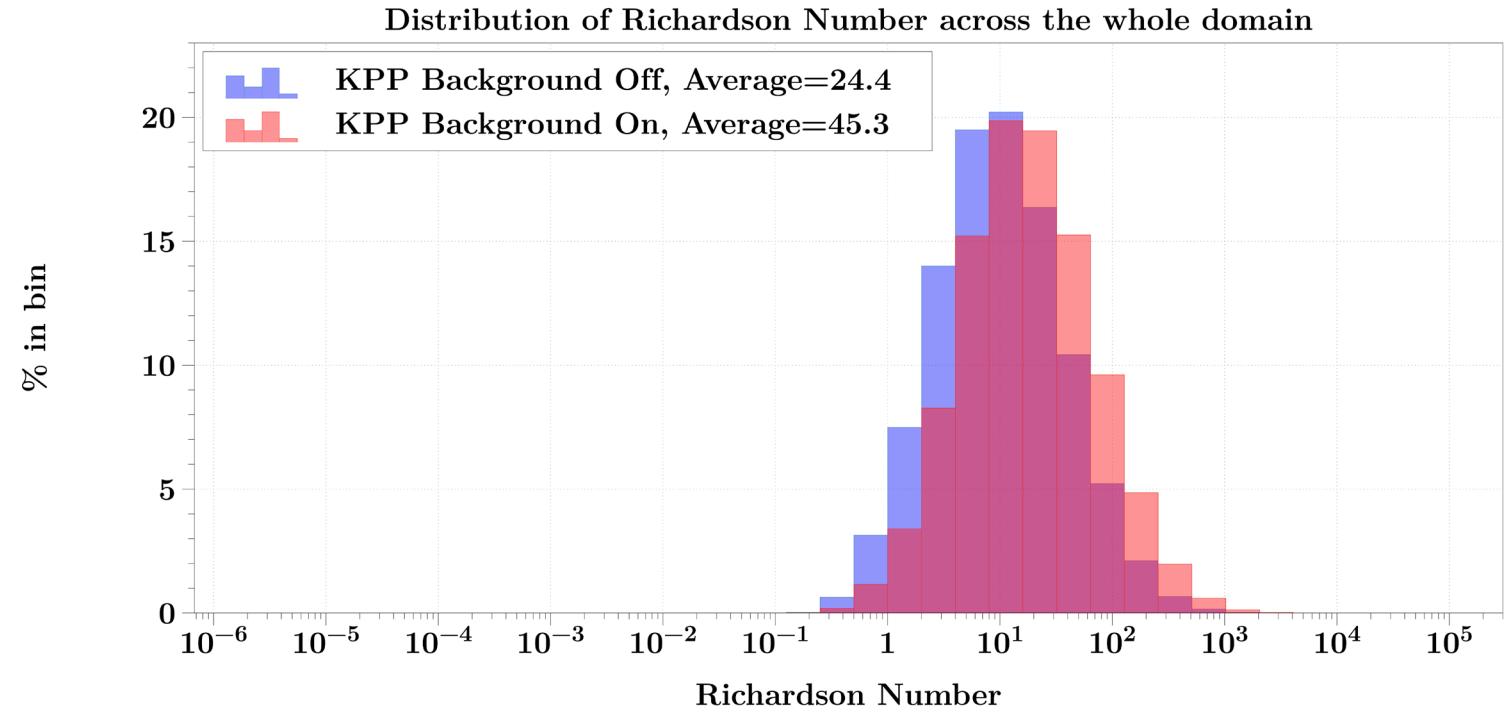


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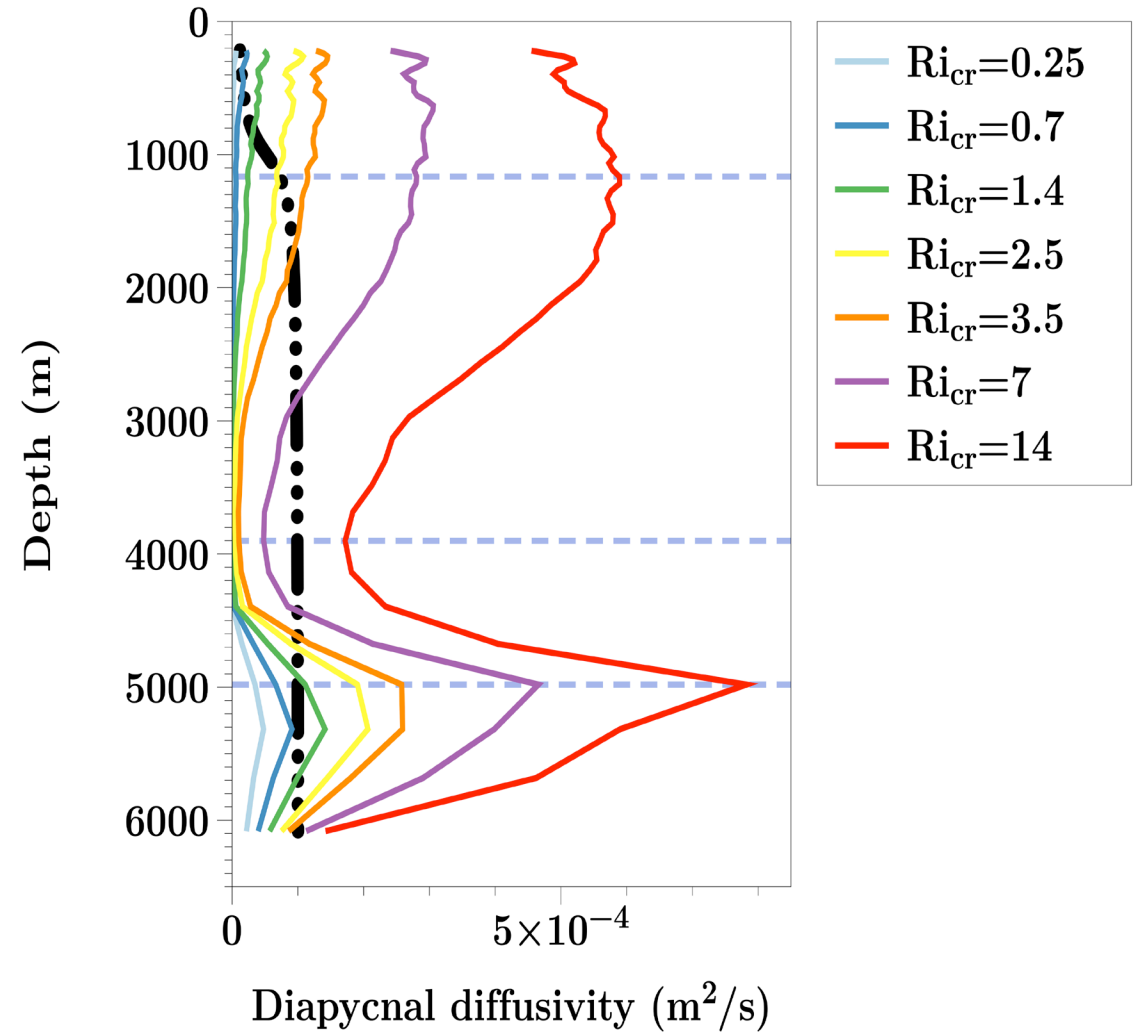
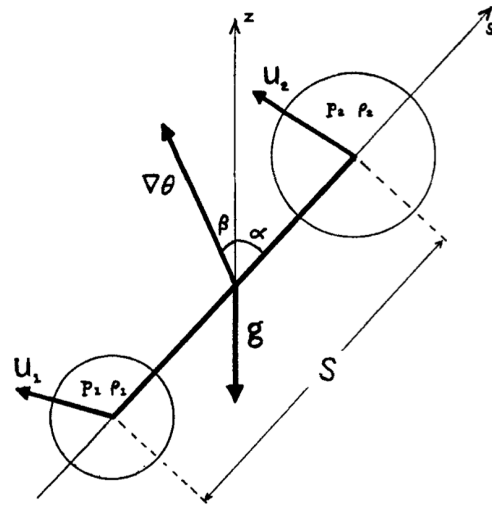


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Left image source: Generalizations of the richardson criterion for the onset of atmospheric turbulence. Quarterly Journal of the Royal Meteorological Society, 97(414), 429–439.

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